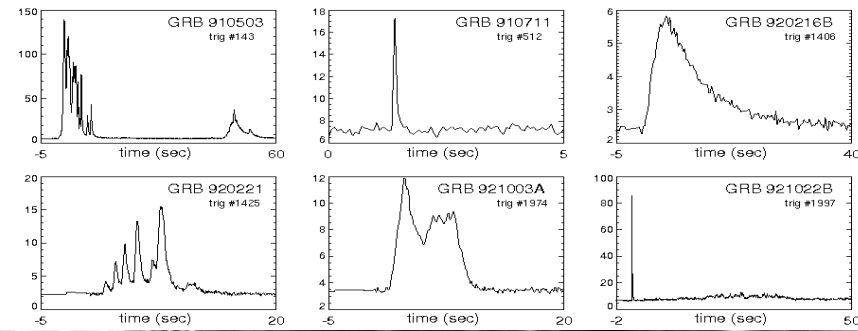




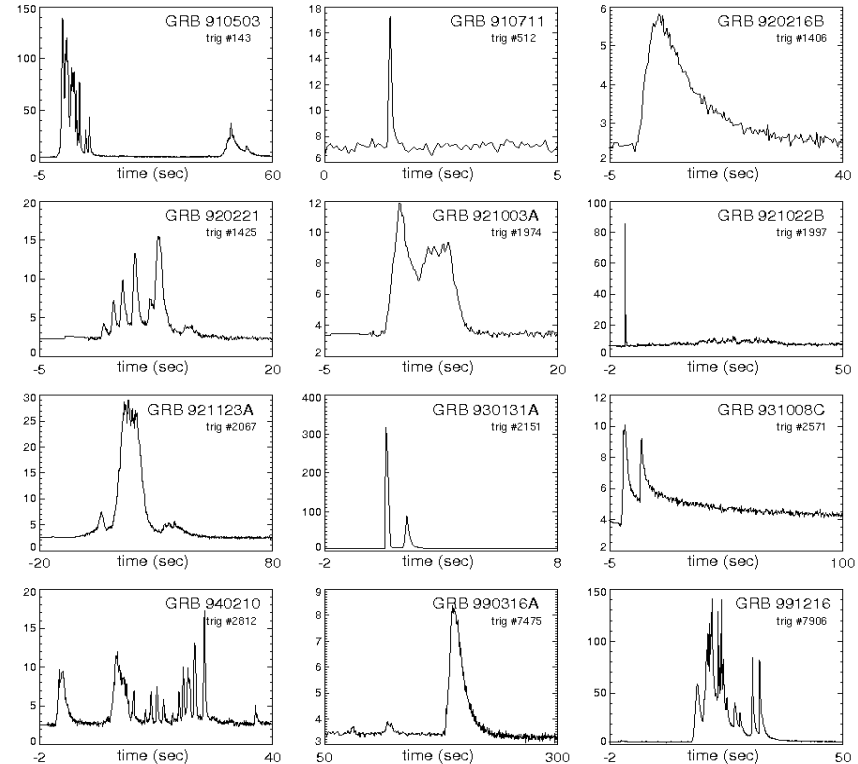
Hypercritical accretion onto neutron stars and the induced gravitational collapse paradigm of gamma-ray bursts associated with supernovae

Jorge Armando Rueda Hernández – ICRANet and Sapienza University of Rome

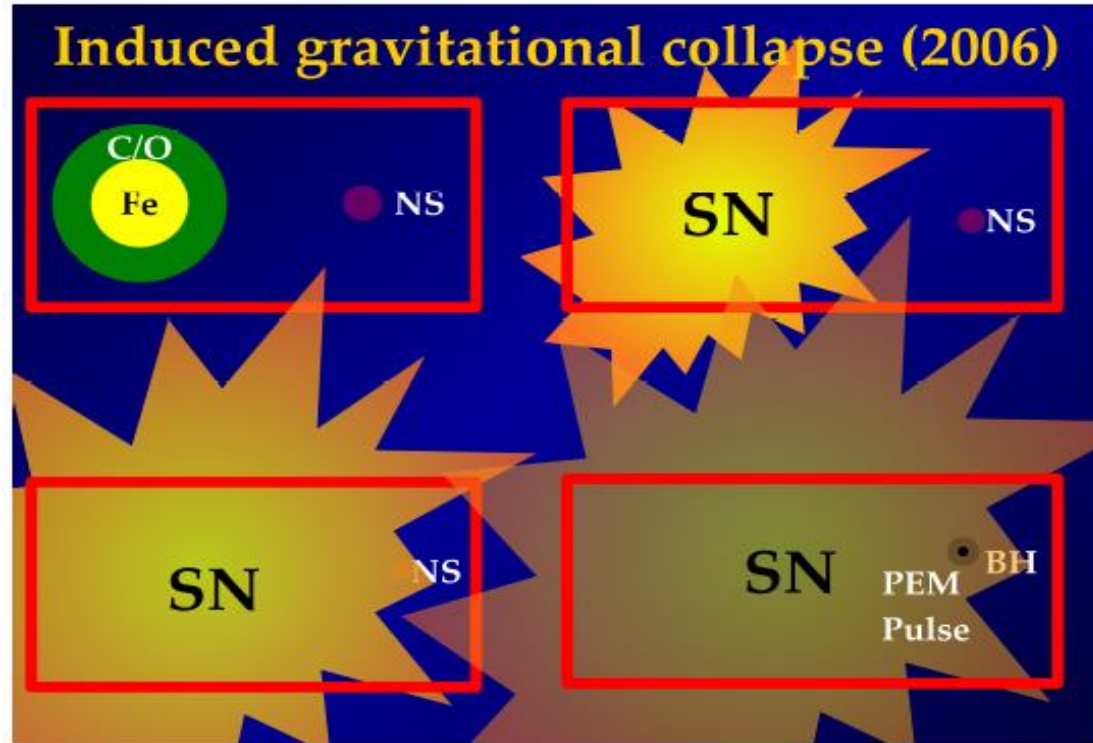
The 1st Sandoval Vallarta Meeting on Relativistic Astrophysics, 30 November-4 December, Mexico City, Mexico
2015

Gamma-Ray Bursts

- GRBs are cosmological explosions (observed up to $z=9.4$ GRB 090429B)
- Most energetic objects (up to a few 10^{54} erg of isotropic energy)
- Complex light-curves but in general characterized by a prompt and an extended afterglow emission
- Duration: “Short” GRBs <2 seconds and “Long” GRBs >2 seconds
- Probe the Physics of *Gravitational Collapse and Black Hole formation*



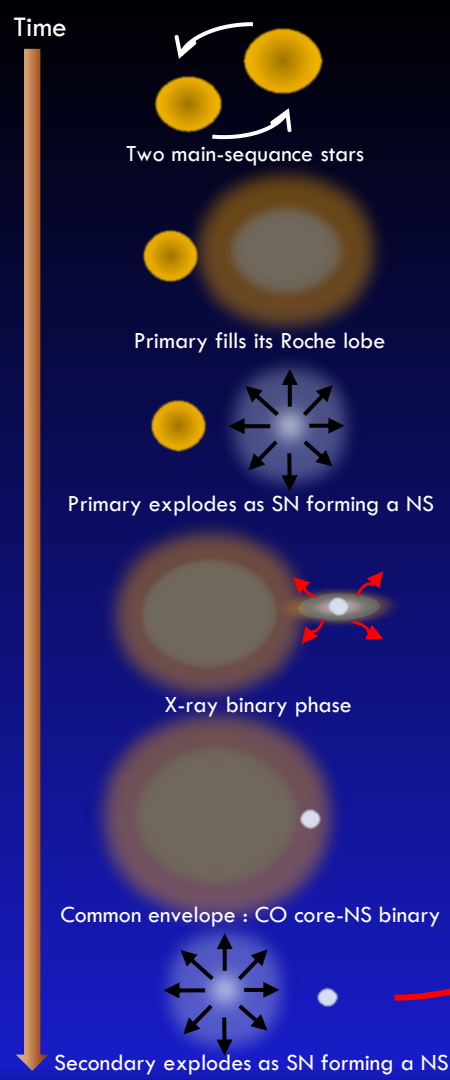
Induced Gravitational Collapse (IGC) Paradigm



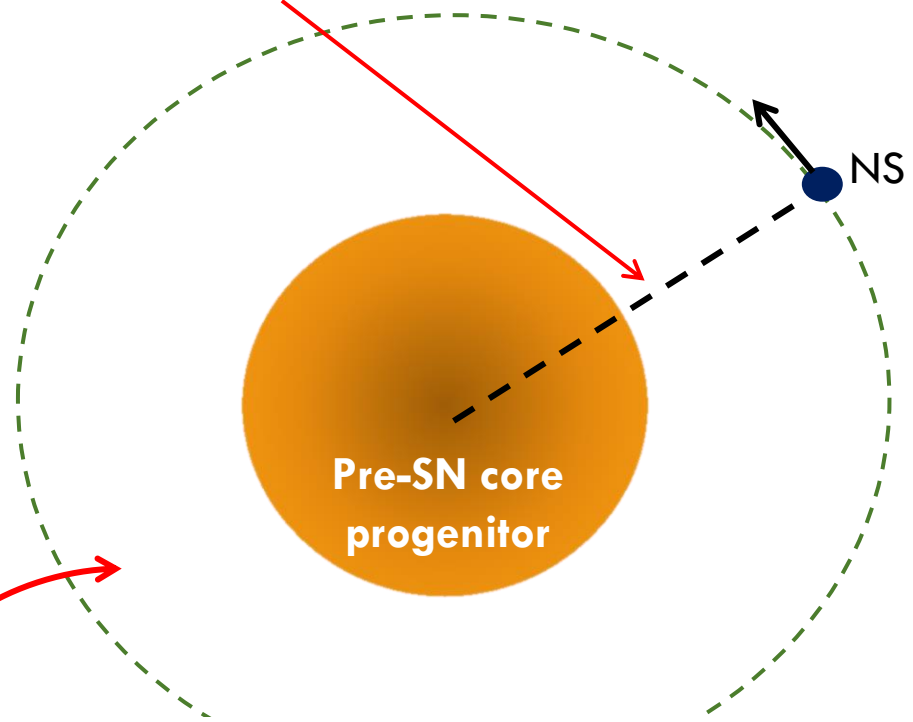
Ruffini et al. MG11-Berlin (2006)

The Induced Gravitational Collapse

Nomoto et al.
1994
Iwamoto et al.
1994
Nomoto et al.
1995



$$v_{\text{orb}} = \sqrt{\frac{G(M_{\text{SN-prog}} + M_{\text{NS}})}{a}} = 1.15 \times 10^8 \left(\frac{M_{\text{SN-prog}} + M_{\text{NS}}}{M_{\odot}} \right)^{1/2} \text{ cm s}^{-1}$$



$$P = \sqrt{\frac{4\pi^2 a^3}{G(M_{\text{SN-prog}} + M_{\text{NS}})}} = 545 \left(\frac{M_{\text{SN-prog}} + M_{\text{NS}}}{M_{\odot}} \right)^{-1/2} \text{ s}$$

Rueda & Ruffini, ApJL (2012)

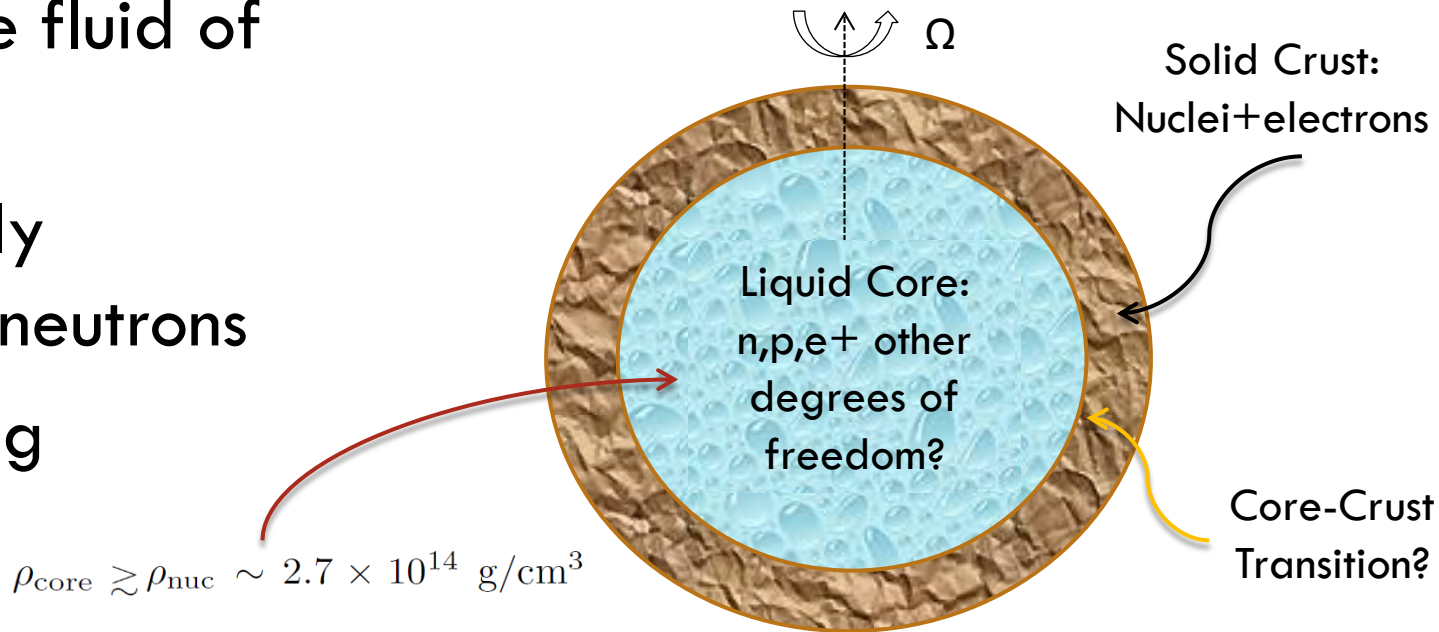
Izzo, Rueda, Ruffini, A&AL (2012)

Oppenheimer & Volkoff (1939) versus Today Neutron Stars

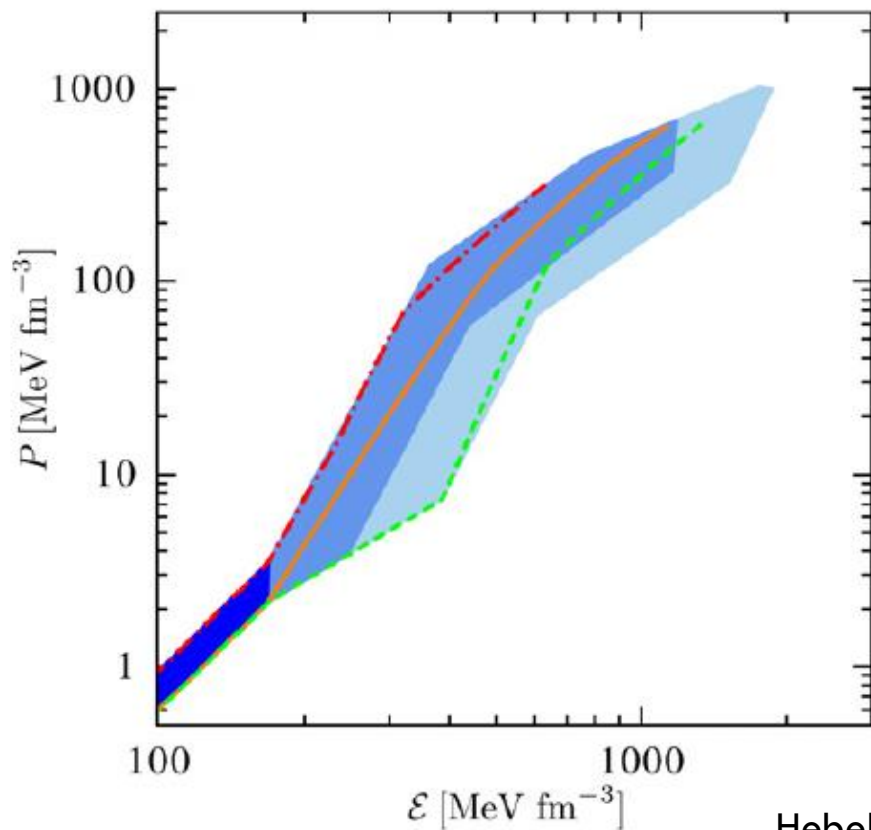
Oppenheimer-Volkoff

- Degenerate fluid of neutrons
- Non-strongly interacting neutrons
- Non-rotating

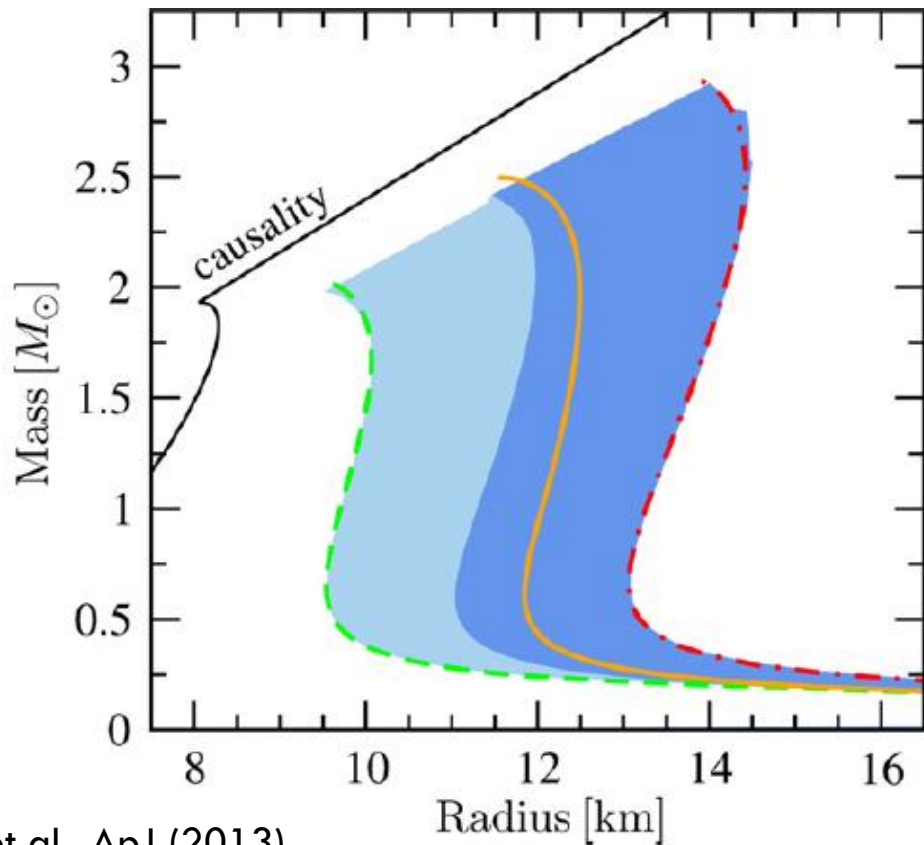
Realistic Neutron Stars



Constraining the nuclear EOS and Mass-Radius Relation



Hebeler et al., ApJ (2013)



NS EOS (Relativistic Mean-Field-RMF- Models)

(Rueda, Ruffini, Xue, Nucl. Phys. A 872, 286, 2011)

$$\mathcal{L} = \mathcal{L}_g + \mathcal{L}_f + \mathcal{L}_\sigma + \mathcal{L}_\omega + \mathcal{L}_\rho + \mathcal{L}_\gamma + \mathcal{L}_{\text{int}}$$

$$\mathcal{L}_g = -\frac{R}{16\pi G},$$

$$\mathcal{L}_\gamma = -\frac{1}{16\pi} F_{\mu\nu} F^{\mu\nu},$$

$$\mathcal{L}_\sigma = \frac{1}{2} \nabla_\mu \sigma \nabla^\mu \sigma - U(\sigma), \quad U(\sigma) = U_0 + U(\sigma, 4)$$

$$\mathcal{L}_\omega = -\frac{1}{4} \Omega_{\mu\nu} \Omega^{\mu\nu} + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu,$$

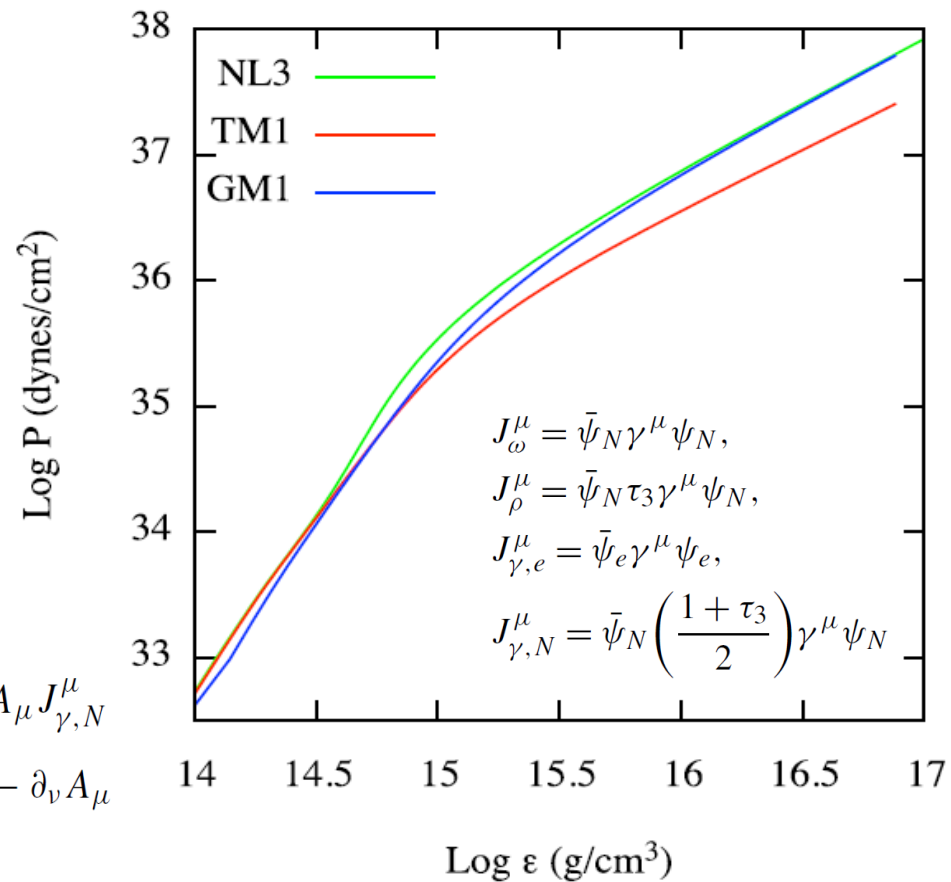
$$\mathcal{L}_\rho = -\frac{1}{4} \mathcal{R}_{\mu\nu} \mathcal{R}^{\mu\nu} + \frac{1}{2} m_\rho^2 \rho_\mu \rho^\mu,$$

$$\mathcal{L}_{\text{int}} = -g_\sigma \sigma \bar{\psi}_N \psi_N - g_\omega \omega_\mu J^\mu_\omega - g_\rho \rho_\mu J^\mu_\rho + e A_\mu J^\mu_{\gamma,e} - e A_\mu J^\mu_{\gamma,N}$$

$$\Omega_{\mu\nu} \equiv \partial_\mu \omega_\nu - \partial_\nu \omega_\mu, \quad \mathcal{R}_{\mu\nu} \equiv \partial_\mu \rho_\nu - \partial_\nu \rho_\mu, \quad F_{\mu\nu} \equiv \partial_\mu A_\nu - \partial_\nu A_\mu$$

$$U_0 \equiv \frac{1}{2} m_\sigma^2 \sigma^2,$$

$$U(\sigma, 4) \equiv \frac{1}{3} g_2 \sigma^3 + \frac{1}{4} g_3 \sigma^4$$



Rotating NS configurations: full rotation in GR

(Cipolletta, Cherubini, Filippi, Rueda, Ruffini, PRD 92, 023007 (2015); arXiv: 1506.05926)

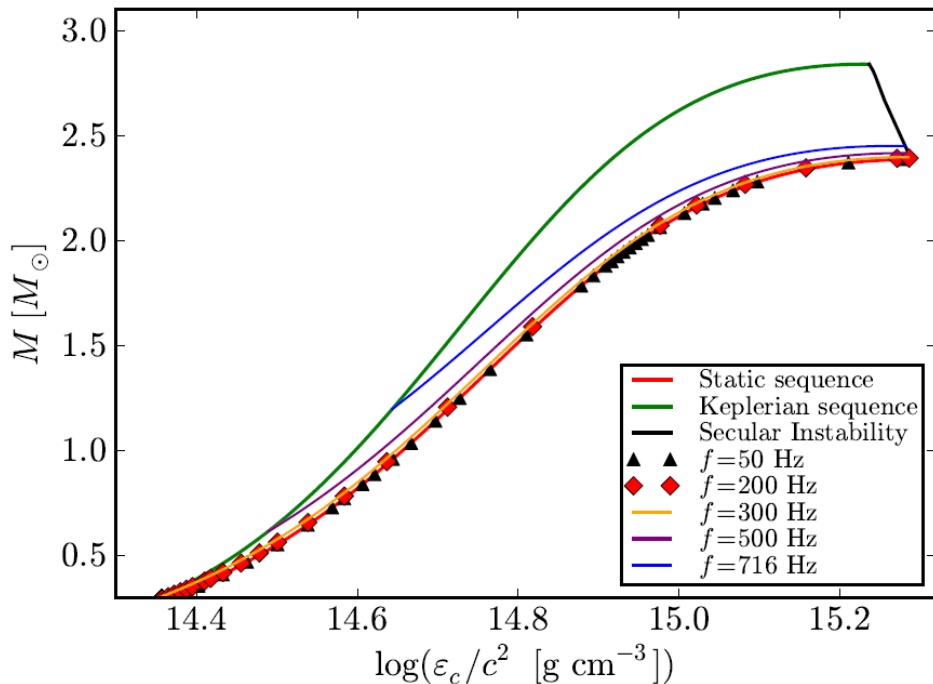
$$ds^2 = -e^{2\nu} dt^2 + e^{2\psi} (d\phi - \omega dt)^2 + e^{2\lambda} (dr^2 + r^2 d\theta^2) \quad T^{\alpha\beta} = (\varepsilon + P) u^\alpha u^\beta + P g^{\alpha\beta}$$

$$\nabla \cdot (B \nabla \nu) = \frac{1}{2} r^2 \sin^2 \theta B^3 e^{-4\nu} \nabla \omega \cdot \nabla \omega + 4\pi B e^{2\zeta - 2\nu} \left[\frac{(\varepsilon + P)(1 + v^2)}{1 - v^2} + 2P \right]$$

$$\nabla \cdot (r^2 \sin^2 \theta B^3 e^{-4\nu} \nabla \omega) = -16\pi r \sin \theta B^2 e^{2\zeta - 4\nu} \frac{(\varepsilon + P)v}{1 - v^2} \quad \nabla \cdot (r \sin(\theta) \nabla B) = 16\pi r \sin \theta B e^{2\zeta - 2\nu} P,$$

$$\begin{aligned} \zeta_{,\mu} = & - \left\{ (1 - \mu^2) \left(1 + r \frac{B_{,r}}{B} \right)^2 + \left[\mu - (1 - \mu^2) \frac{B_{,r}}{B} \right]^2 \right\}^{-1} \left[\frac{1}{2} B^{-1} \left\{ r^2 B_{,rr} - [(1 - \mu^2) B_{,\mu}]_{,\mu} - 2\mu B_{,\mu} \right\} \right. \\ & \times \left\{ -\mu + (1 - \mu^2) \frac{B_{,\mu}}{B} \right\} + r \frac{B_{,r}}{B} \left[\frac{1}{2} \mu + \mu r \frac{B_{,r}}{B} + \frac{1}{2} (1 - \mu^2) \frac{B_{,\mu}}{B} \right] + \frac{3}{2} \frac{B_{,\mu}}{B} \left[-\mu^2 + \mu (1 - \mu^2) \frac{B_{,\mu}}{B} \right] \\ & - (1 - \mu^2) r \frac{B_{,\mu r}}{B} \left(1 + r \frac{B_{,r}}{B} \right) - \mu r^2 (\nu_{,r})^2 - 2(1 - \mu^2) r \nu_{,\mu} \nu_{,r} + \mu (1 - \mu^2) (\nu_{,\mu})^2 - 2(1 - \mu^2) r^2 B^{-1} B_{,r} \nu_{,\mu} \nu_{,r} \\ & + (1 - \mu^2) B^{-1} B_{,\mu} \left[r^2 (\nu_{,r})^2 - (1 - \mu^2) (\nu_{,\mu})^2 \right] + (1 - \mu^2) B^2 e^{-4\nu} \left\{ \frac{1}{4} \mu r^4 (\omega_{,r})^2 + \frac{1}{2} (1 - \mu^2) r^3 \omega_{,\mu} \omega_{,r} \right. \\ & \left. \left. - \frac{1}{4} \mu (1 - \mu^2) r^2 (\omega_{,\mu})^2 + \frac{1}{2} (1 - \mu^2) r^4 B^{-1} B_{,r} \omega_{,\mu} \omega_{,r} - \frac{1}{4} (1 - \mu^2) r^2 B^{-1} B_{,\mu} \left[r^2 (\omega_{,r})^2 - (\mu^2) (\omega_{,\mu})^2 \right] \right\} \right] \end{aligned}$$

Rotating NS configurations: secular instability line

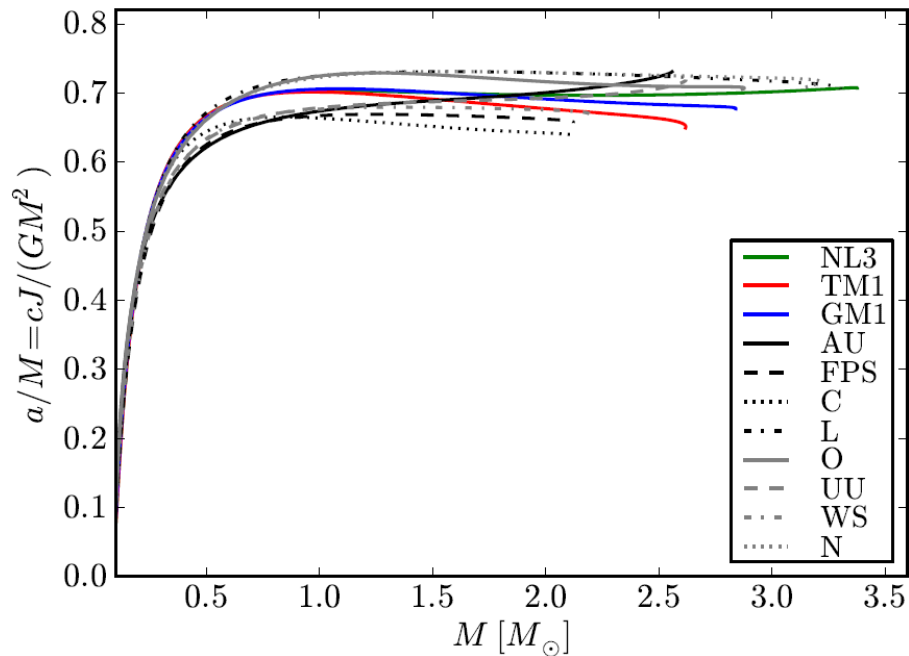
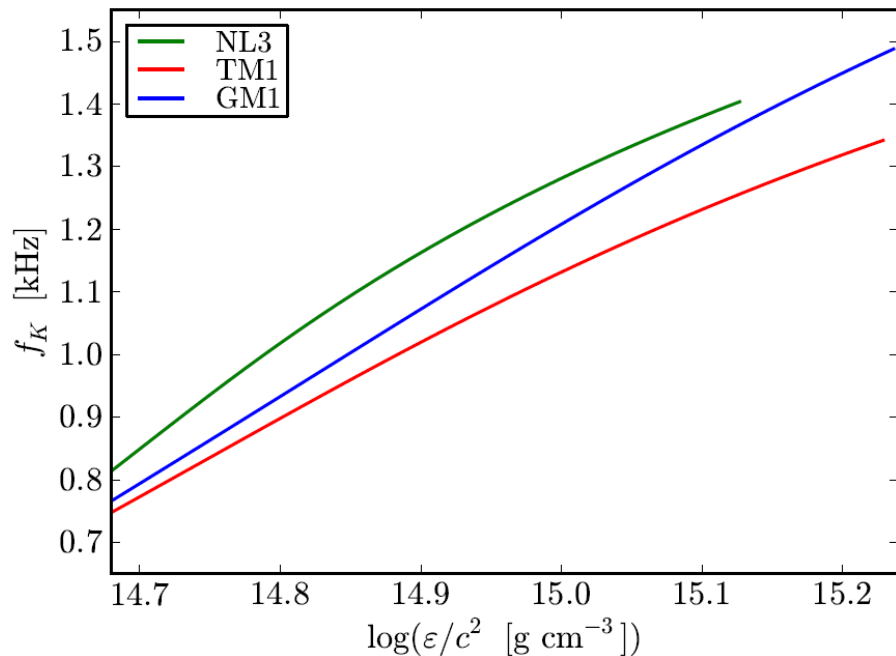


$$M_{\text{NS}}^{\text{crit}} = M_{\text{NS}}^{J=0} (1 + k j_{\text{NS}}^p)$$

	$M_{\text{crit}}^{J=0}$ $M_\odot$	$R_{\text{crit}}^{J=0}$ km	$M_{\text{max}}^{J \neq 0}$ $M_\odot$	$R_{\text{max}}^{J \neq 0}$ km	f_K kHz	p	k
NL3	2.81	13.49	3.38	17.35	1.34	1.68	0.006
GM1	2.39	12.56	2.84	16.12	1.49	1.69	0.011
TM1	2.20	12.07	2.62	15.98	1.40	1.61	0.017

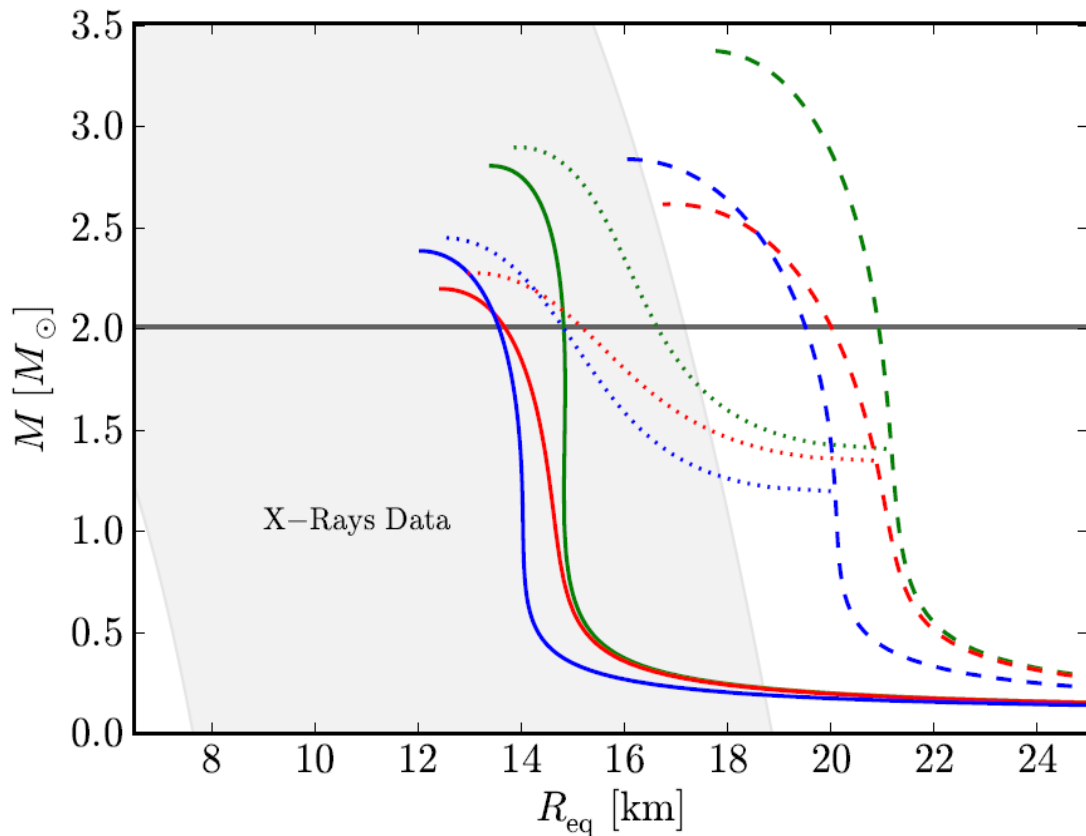
Taken from Cipolletta, et al. PRD 92, 023007 (2015)
arXiv: 1506.05926

Rotating NS configurations: full rotation in GR



Figures from Cipolletta, et al. PRD 92, 023007 (2015); arXiv: 1506.05926

NS Mass-Radius Relation: Observational Constraints

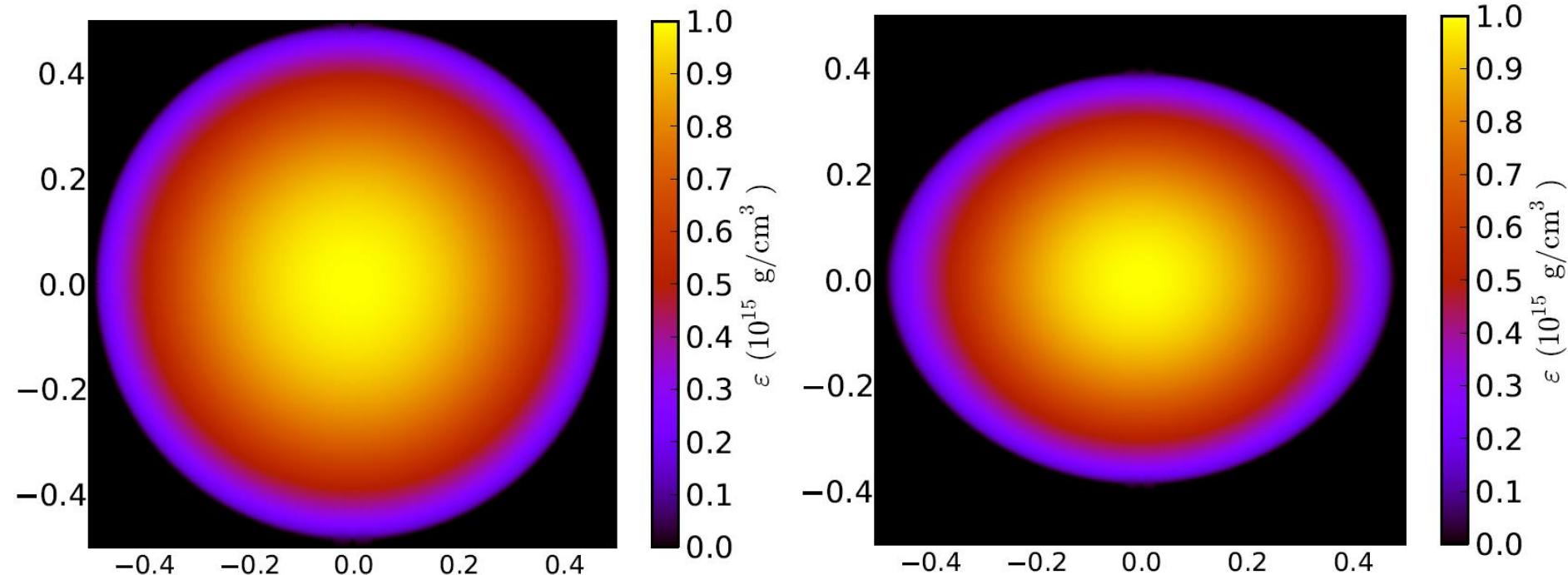


- Maximum NS mass observed:
2 M_{sun}
(Antoniadis et al., Science (2013))
- Fastest NS observed: $f=716$ Hz
(Demorest et al., Science (2006))
- Radii from X-ray emission: mainly from low-mass X-ray binaries (LMXBs), and X-ray isolated NSs (XINSs): shaded area

(Lattimer & Steiner, EPJ (2014))

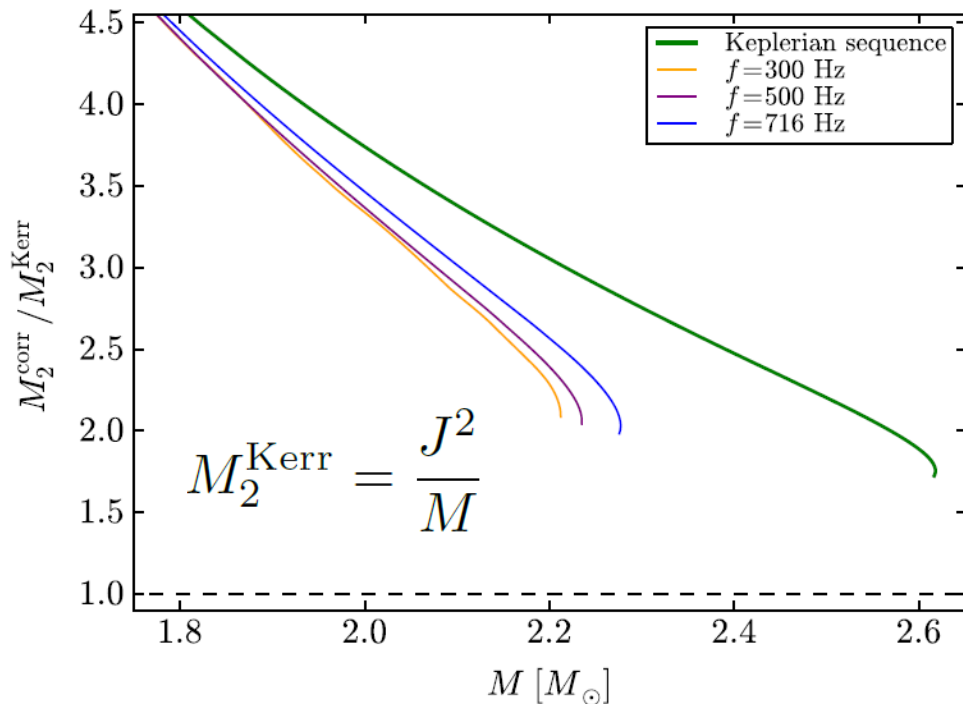
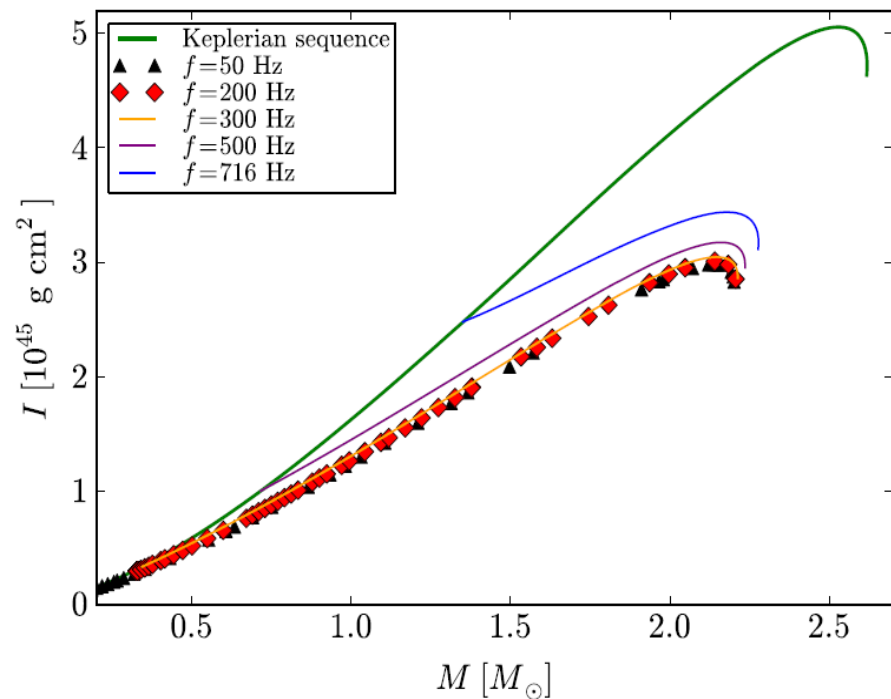
Rotating NS: Deformation

(Cipolletta, Cherubini, Filippi, Rueda, Ruffini, PRD 92, 023007 (2015); arXiv: 1506.05926)



NS moment of inertia and quadrupole moment

(Cipolletta, Cherubini, Filippi, Rueda, Ruffini, PRD 92, 023007 (2015); arXiv: 1506.05926)



Neutron Star Binding Energy

(Cipolletta, Cherubini, Filippi, Rueda, Ruffini, PRD 2015; arXiv: 1506.05926)

Static Configurations

$$\frac{M_b}{M_\odot} \approx \frac{M}{M_\odot} + \frac{13}{200} \left(\frac{M}{M_\odot} \right)^2$$

$c \text{ J}/(G M_\text{sun}^2)$

Rotating Configurations

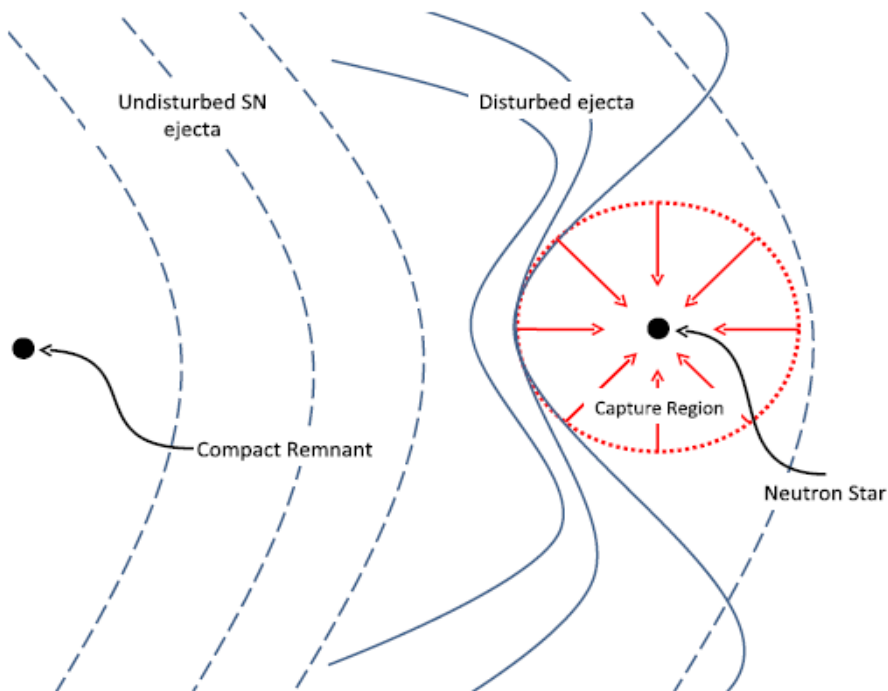
$$\frac{M_b}{M_\odot} = \frac{M}{M_\odot} + \frac{13}{200} \left(\frac{M}{M_\odot} \right)^2 \left(1 - \frac{1}{130} j^{1.7} \right)$$

First estimates of the accretion process

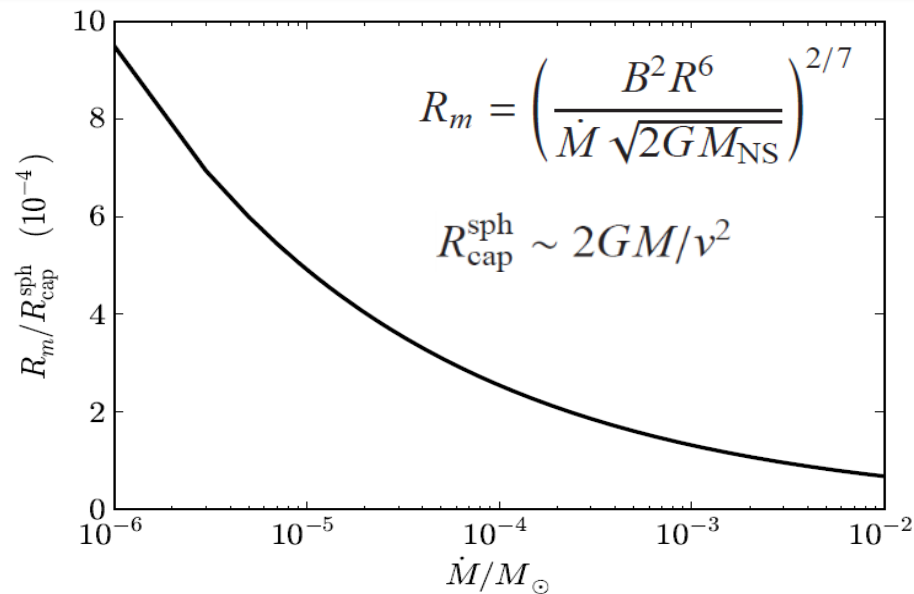
Rueda & Ruffini, ApJ Lett. 758, L7 (2012)

Izzo, Rueda, Ruffini, A&A Lett. 548, L5 (2012)

$$R_{\text{cap}} = \frac{2GM_{\text{NS}}}{v_{\text{rel,ej}}^2}, \quad v_{\text{rel,ej}} = \sqrt{v_{\text{orb}}^2 + v_{\text{ej}}^2}, \quad r_{\text{ej}} = \sigma t^n$$



The relative importance of the magnetosphere



See also Toropina, Romanova, Lovelace, MNRAS 420, 810 (2012)

How long is the accretion process?
Does the neutron star reach maximum mass?

**Which is the initial mass of the
neutron star?**

**Which is the maximum stable
mass of a neutron star?**

Improvements to the First IGC Scenario

- **SN core density and SN initial velocity profiles from numerical simulations**
- **SN core and NS masses from binary evolution codes**
- **Hydrodynamics inside the Bondi accretion region: photon trapping radius, neutrino emission**
- **Characteristic emission from the accretion process**
 - ...

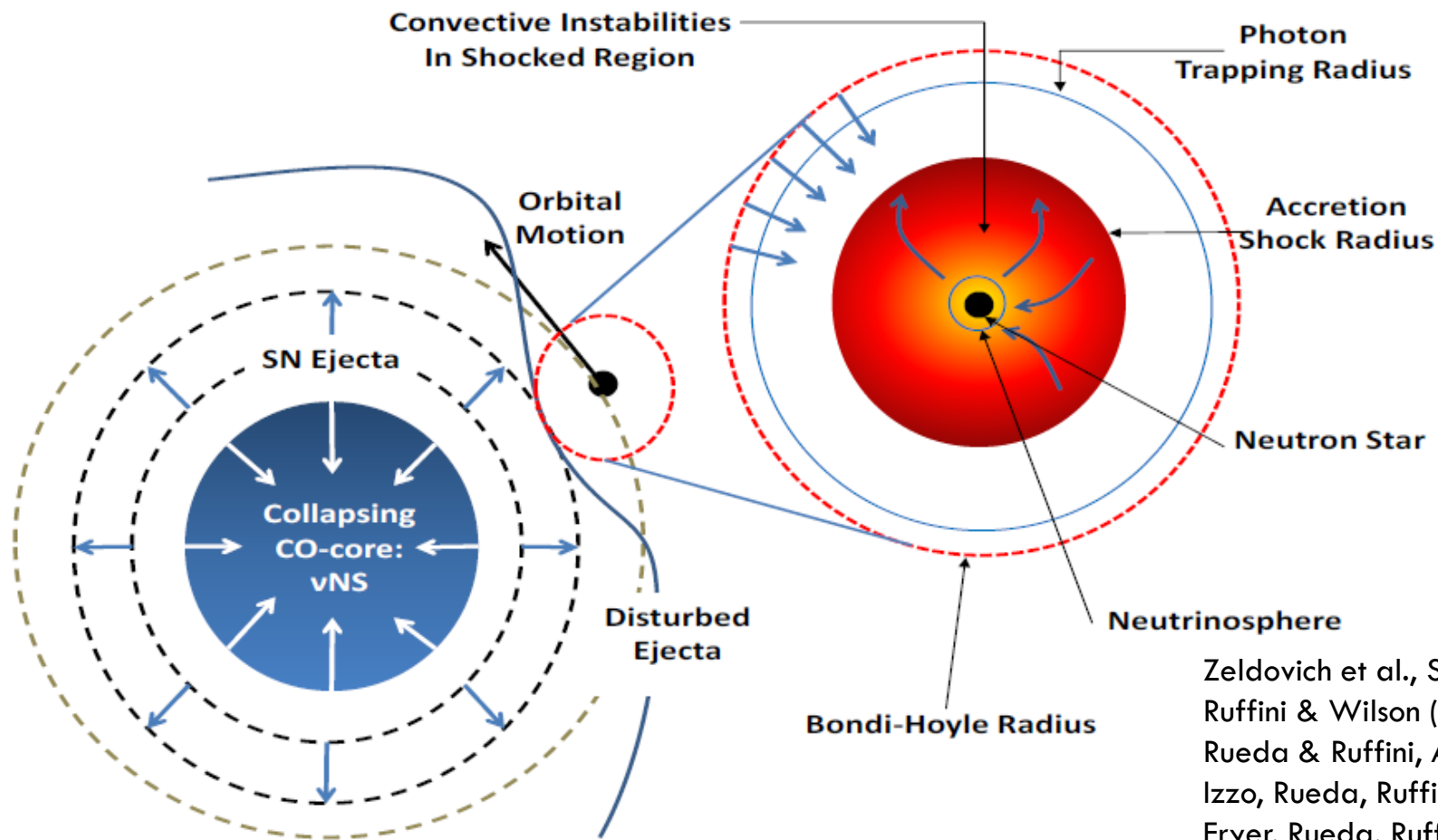
Hypercritical Accretion

Conditions for Eddington limited accretion:

- Potential energy is released in the form of photons
- Inflowing material and outflowing radiation are spherically symmetric
- Photons can flow and deposit momentum to the inflowing material
 - Opacity is dominated by electron scattering

***NONE OF THE ABOVE CONDITIONS IS SATISFIED IN THE
IGC BINARY SYSTEM !!!***

Binary Driven Hypercritical Accretion in the IGC

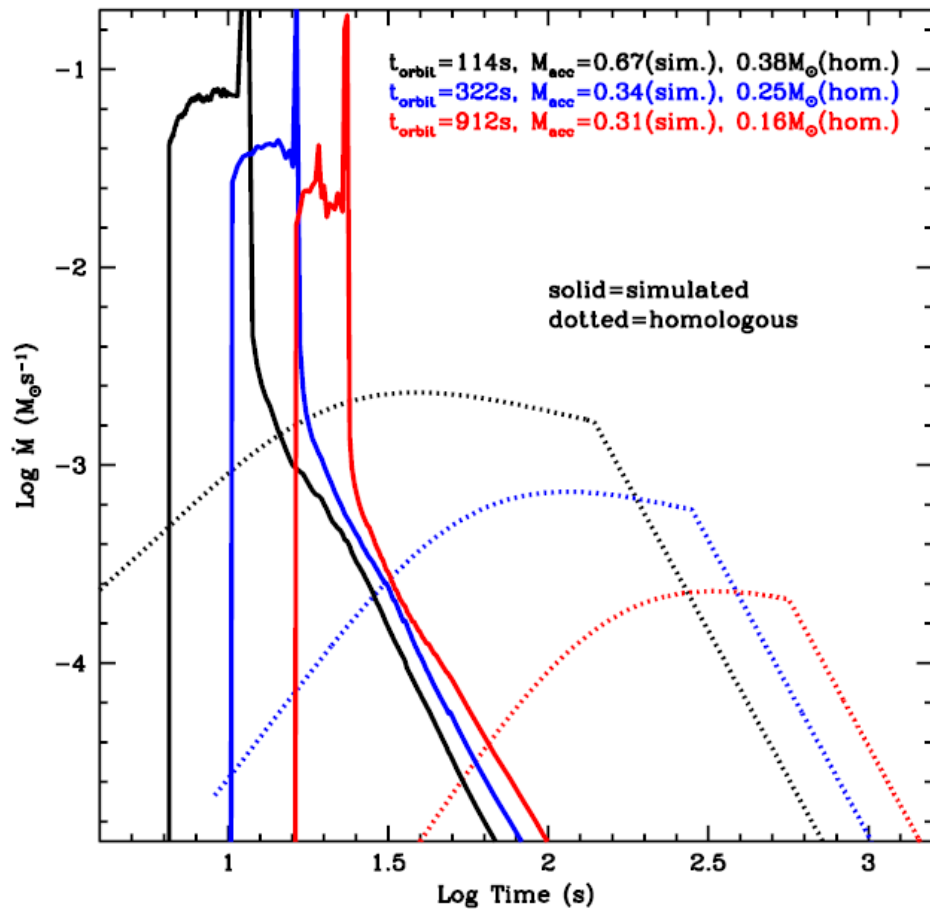
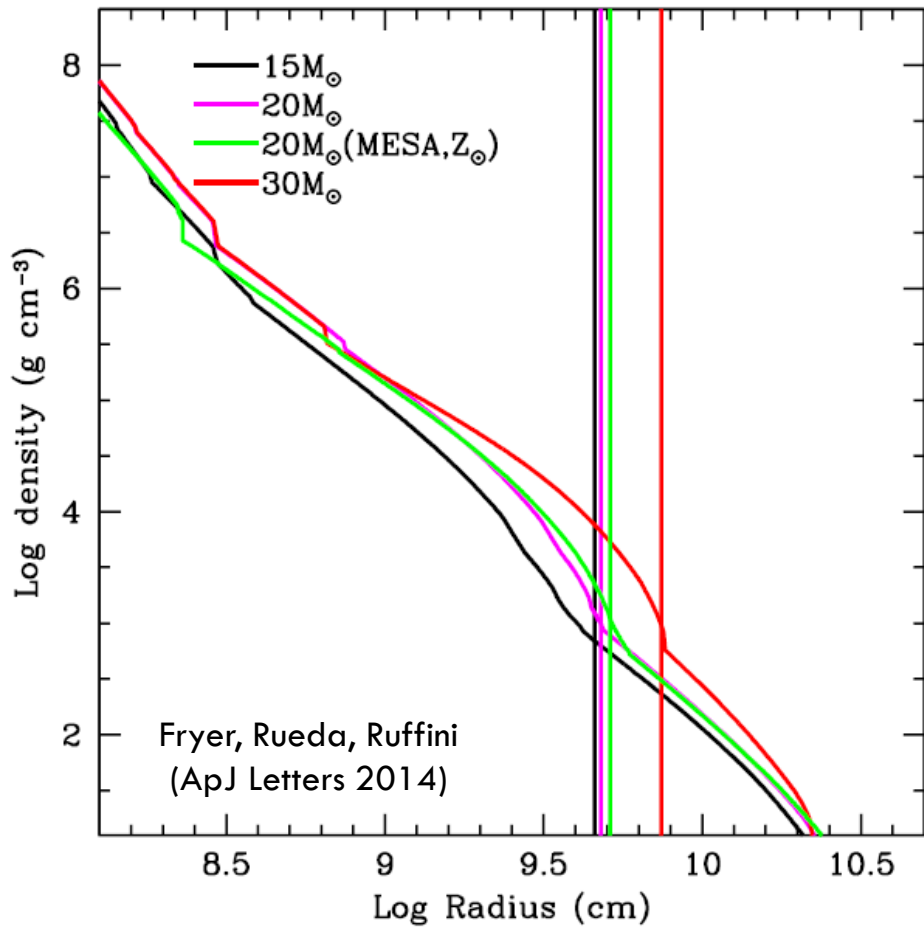


Zeldovich et al., Sov. Astron. (1972)
Ruffini & Wilson (1973)
Rueda & Ruffini, ApJL (2012)
Izzo, Rueda, Ruffini, A&AL (2012)
Fryer, Rueda, Ruffini, ApJL (2014)

$$\dot{M}_{\text{BHL}} = 4\pi r_{\text{BHL}}^2 \rho (v^2 + c_s^2)^{1/2}$$

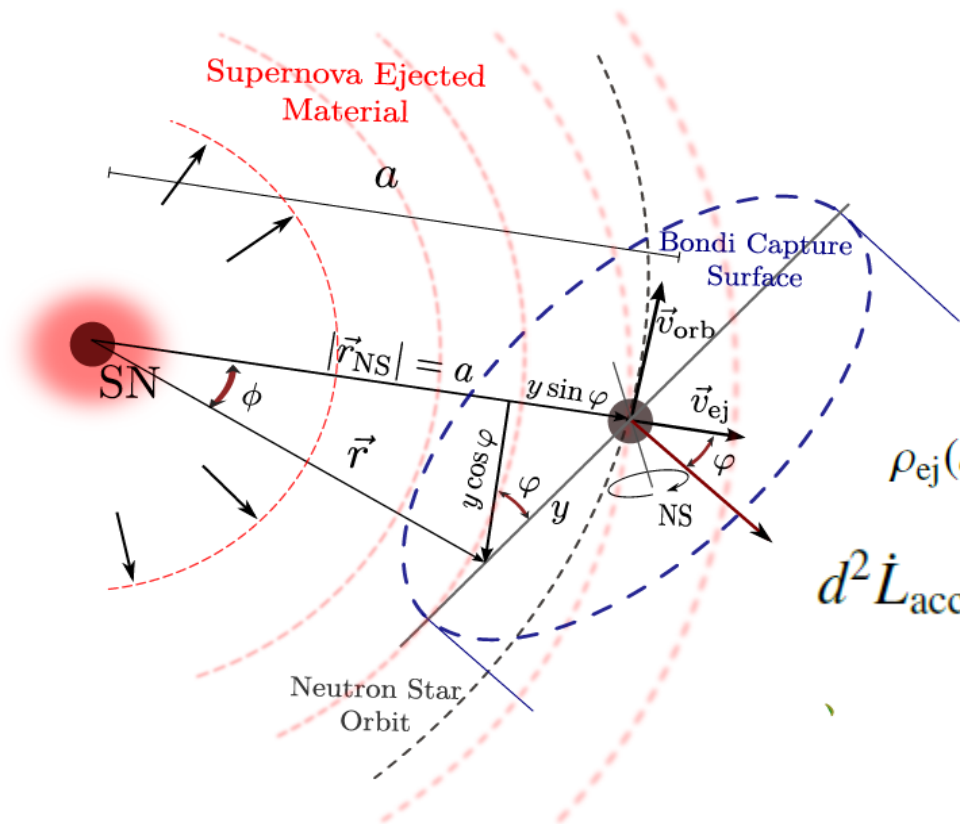
$$r_{\text{BHL}} = \frac{GM_{\text{NS}}}{v^2 + c_s^2}$$

$$r_{\text{trapping}} = \min[(\dot{M}_{\text{BHL}}\kappa)/(4\pi c), r_{\text{BHL}}]$$



On the role of angular momentum in BdHNe

(Becerra, Cipolletta, Fryer, Rueda, Ruffini, ApJ 2015; arXiv: 1505.07580)



$$\dot{M}_B(t) = \pi \rho_{\text{ej}} R_{\text{cap}}^2 \sqrt{v_{\text{rel}}^2 + c_{\text{s,ej}}^2}$$

$$R_{\text{cap}}(t) = \frac{2GM_{\text{NS}}(t)}{v_{\text{rel}}^2 + c_{\text{s,ej}}^2}$$

$$\rho_{\text{ej}}(a) \simeq \rho_{\text{ej}}(a)(1 + \epsilon_{\rho}y) \quad \text{and} \quad v_{\text{rel}}(a) \simeq v_{\text{rel}}(a)(1 + \epsilon_{\nu}y)$$

$$d^2 \dot{L}_{\text{acc}} = \rho_{\text{ej}}(a) v_{\text{rel}}^2(a) \left[y + (\epsilon_{\rho} + 2\epsilon_{\nu})y^2 \right] dy dz$$

$$y^2 + z^2 = R_{\text{cap}}^2 = \left(\frac{2GM_{\text{NS}}}{v_{\text{rel}}^2(a, t)} \right)^2 (1 - 4\epsilon_{\nu}y)$$

On the role of angular momentum in BdHNe

(Becerra, Cipolletta, Fryer, Rueda, Ruffini, ApJ 2015; arXiv:1505.07580)

$$\dot{L}_{\text{acc}} = \frac{\pi}{2} \left(\frac{1}{2} \epsilon_\rho - 3 \epsilon_\nu \right) \rho_{\text{ej}}(a, t) v_{\text{rel}}^2(a, t) R_{\text{cap}}^4(a, t)$$

$$\rho_{\text{ej}}(x, t) = \rho_{\text{ej}}(x, t_0) \frac{M_{\text{env}}(t)}{M_{\text{env}}(t_0)} \left(\frac{R_{0\text{star}}}{R_{\text{star}}(t)} \right)^3 \quad x \equiv \frac{r}{R_{\text{star}}}$$

$$R_{\text{star}}(t) = R_{0\text{star}} \left(\frac{t}{t_0} \right)^n \quad v_{\text{ej}}(r, t) = n \frac{r}{t}$$

$$\rho_{\text{ej}}(r, t) \approx \rho_{\text{ej}}(a, t) \left(1 + \frac{1}{\rho_{\text{ej}}(a, t)} \left. \frac{\partial \rho_{\text{ej}}}{\partial r} \right|_{(a, t)} \delta r \right)$$

$$v_{\text{ej}}(r, t) \approx v_{\text{ej}}(a, t) \left(1 + \frac{1}{v_{\text{ej}}(a, t)} \left. \frac{\partial v_{\text{ej}}}{\partial r} \right|_{(a, t)} \delta r \right)$$

$$\dot{L}_{\text{acc}} = 8\pi \rho_{\text{core}} \left(\frac{R_{\text{core}}}{a} \right)^m \frac{GM_{\text{NS}}(t_0) a^2}{(1+q)^3} H(y)$$

$$H(y) = y^{n(m-3)} \left(\frac{M_{\text{G}}(M_b)}{M_{\text{NS}}(t_0)} \right)^4 (1 - \chi \mu_B) \left(1 + \frac{\eta}{y^2} \right)^{-7/2} \left(\frac{m}{2} + \frac{6\eta}{y^2 + \eta} \right)$$

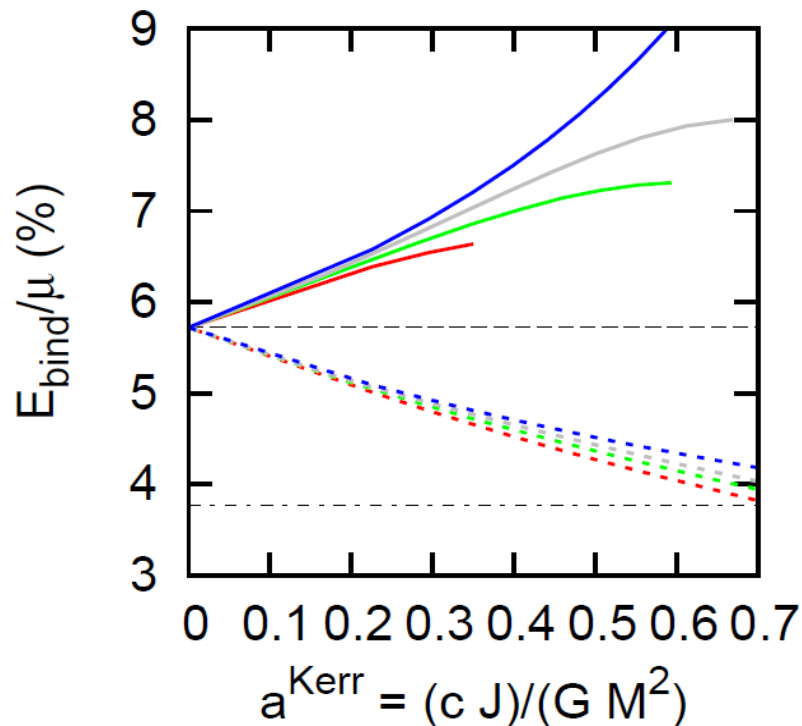
$$\frac{\dot{\mu}_B}{(1 - \chi \mu_B) M_{\text{G}}(M_b)^2} = \frac{t_0}{\tau_B} \frac{y^{n(m-3)}}{\hat{r}^m} \left[1 + \eta \left(\frac{\hat{r}}{y} \right)^2 \right]^{-3/2}$$

$$y \equiv \frac{t}{t_0}, \quad \mu_B \equiv \frac{M_B(y)}{M_{\text{NS}}(t_0)}, \quad \hat{r} \equiv 1 - \frac{R_B}{a}$$

$$\tau_B \equiv \frac{M_{\text{NS}}(t_0) v_{\text{orb}}^3}{4\pi G^2 \rho_{\text{ej}}(a, t_0)}, \quad \chi \equiv \frac{M_{0\text{NS}}}{M_{\text{env}}(t_0)}, \quad \eta \equiv \left(\frac{n}{t_0} \frac{a}{v_{\text{orb}}} \right)^2$$

Mostly bound circular orbit around rotating NSs

(Cipolletta, Rueda, Ruffini, PRD, submitted)



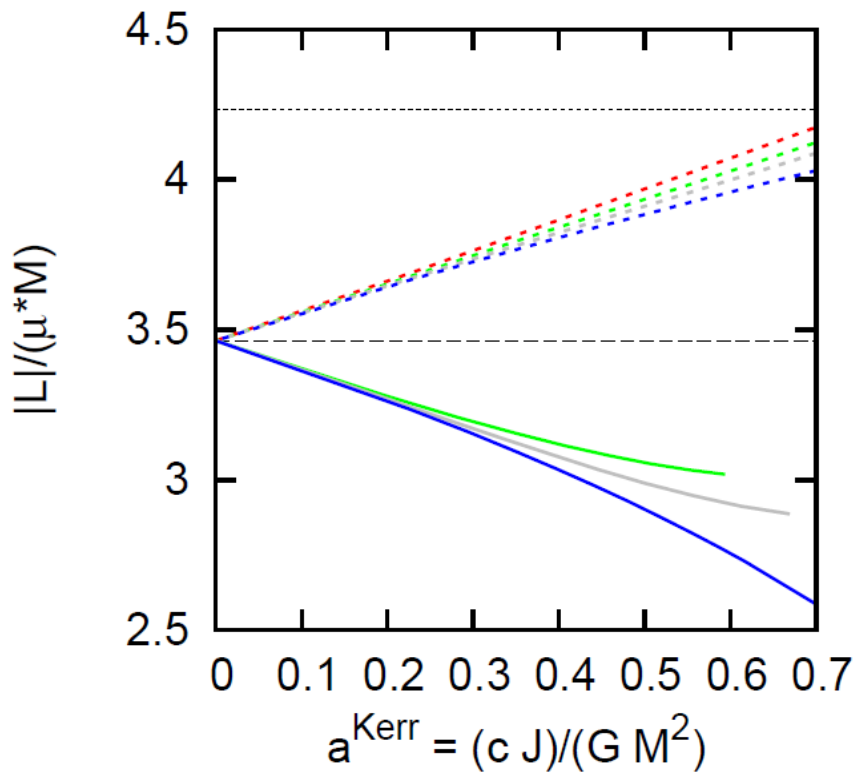
Schw = 5.72 %
Ex Kerr = 3.77 %
 E_+ for $M = 2$
 E_- for $M = 2$
 E_+ for $M = 2.4$
 E_- for $M = 2.4$
 E_+ for $M = 2.8$
 E_- for $M = 2.8$
 E_+ Kerr
 E_- Kerr

Numerical Fit (common for the RMF EOS !):

$$\tilde{E} = 0.9428 - 0.0132 \left(\frac{j}{M} \right)^{0.85}$$

Mostly bound circular orbit around rotating NSs

(Cipolletta, Rueda, Ruffini, PRD, submitted)



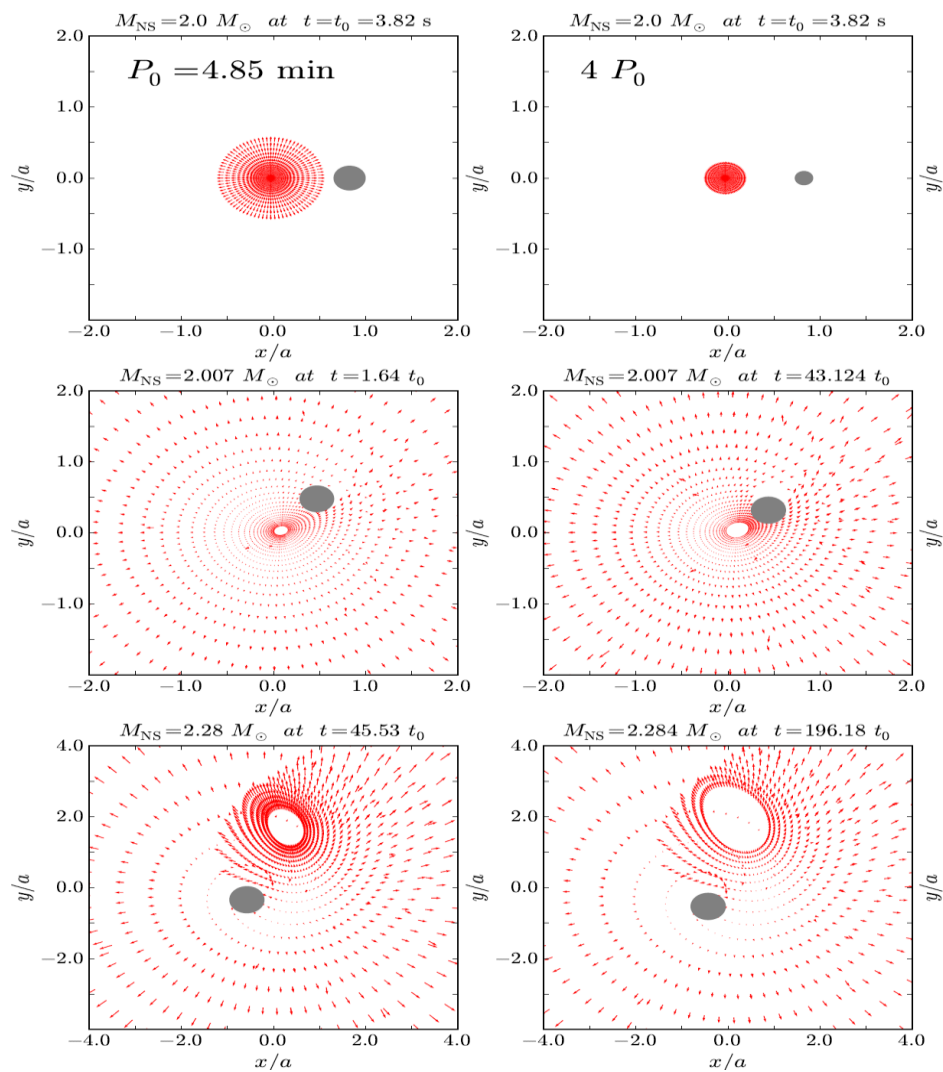
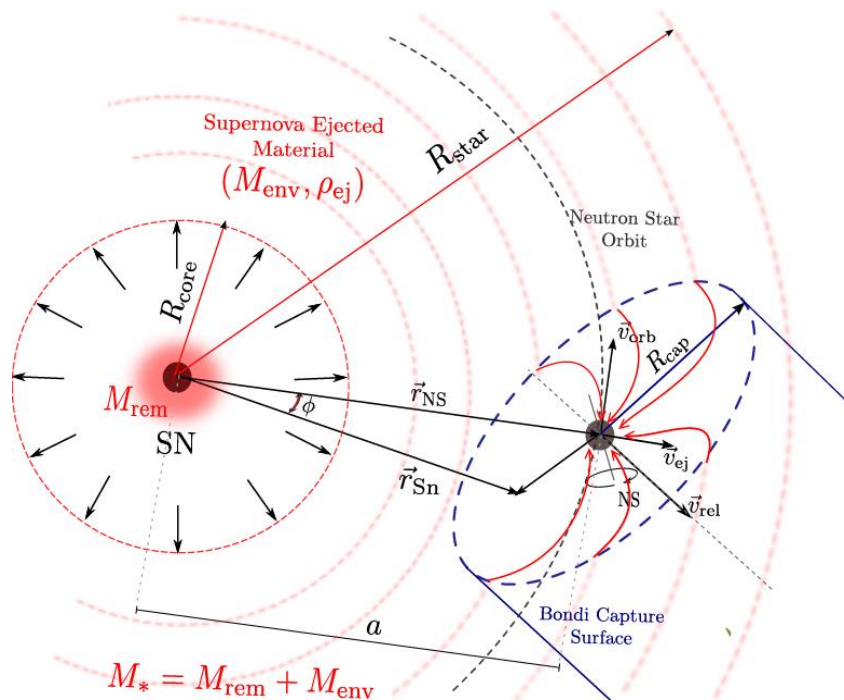
$L_{\pm}$	Schw = $2\sqrt{3}$	---
$L_{\pm}$	ex Kerr = $22/(3\sqrt{3})$	-.-.-
L_{+}	for $M = 2$	—
L_{-}	for $M = 2$	- - -
L_{+}	for $M = 2.4$	—
L_{-}	for $M = 2.4$	- - -
L_{+}	for $M = 2.8$	—
L_{-}	for $M = 2.8$	- - -
L_{+}	Kerr	—
L_{-}	Kerr	- - -

Numerical Fit (common for the RMF EOS !):

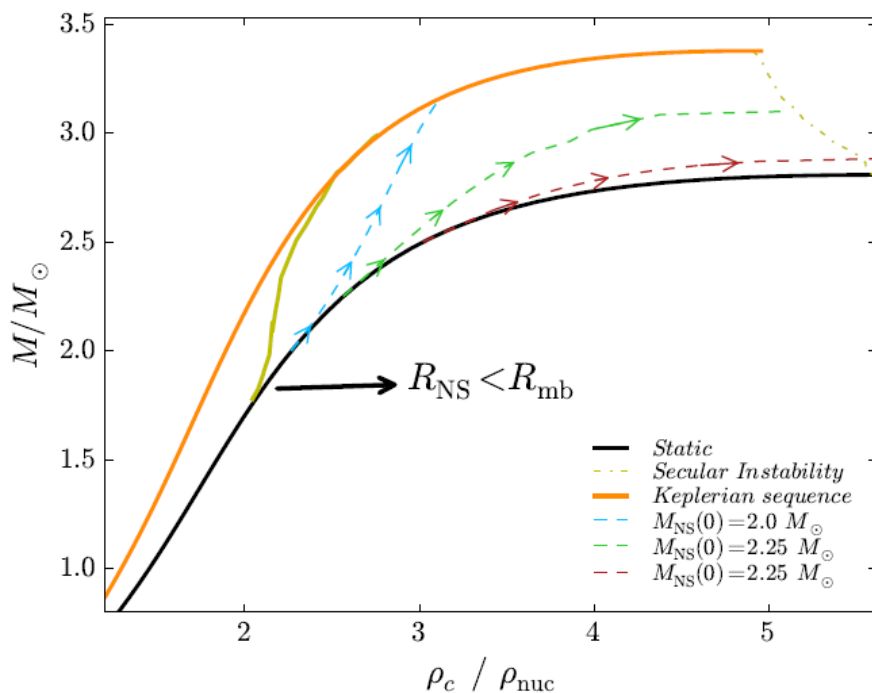
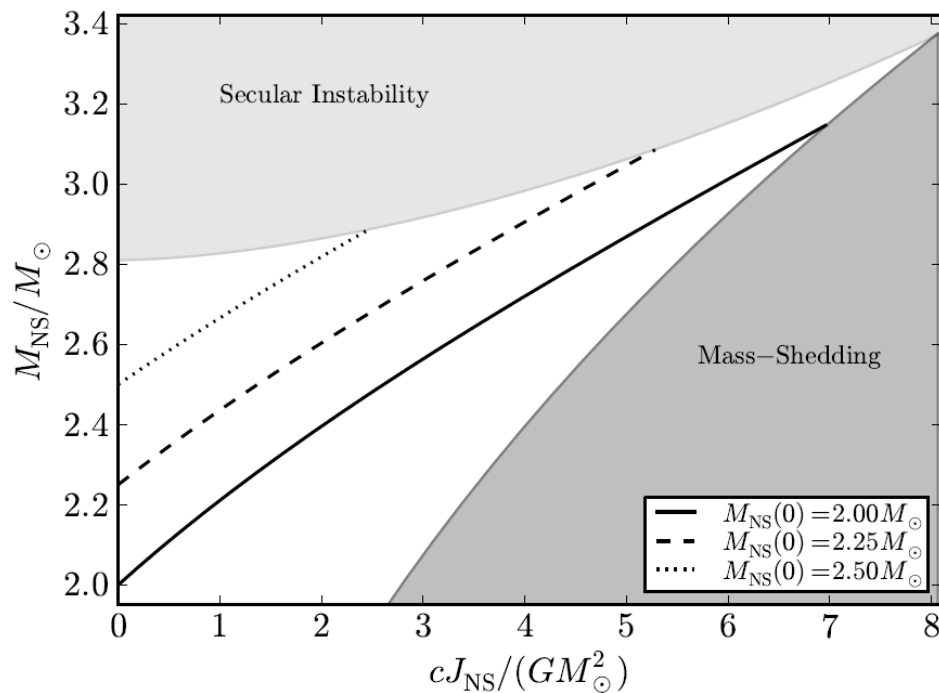
$$\tilde{L} = 3.464 - 0.37 \left(\frac{j}{M} \right)^{0.85}$$

NS evolution during hypercritical accretion

Becerra, Cipolletta, Fryer, Rueda, Ruffini, ApJ 2015:
arXiv:1505.07580

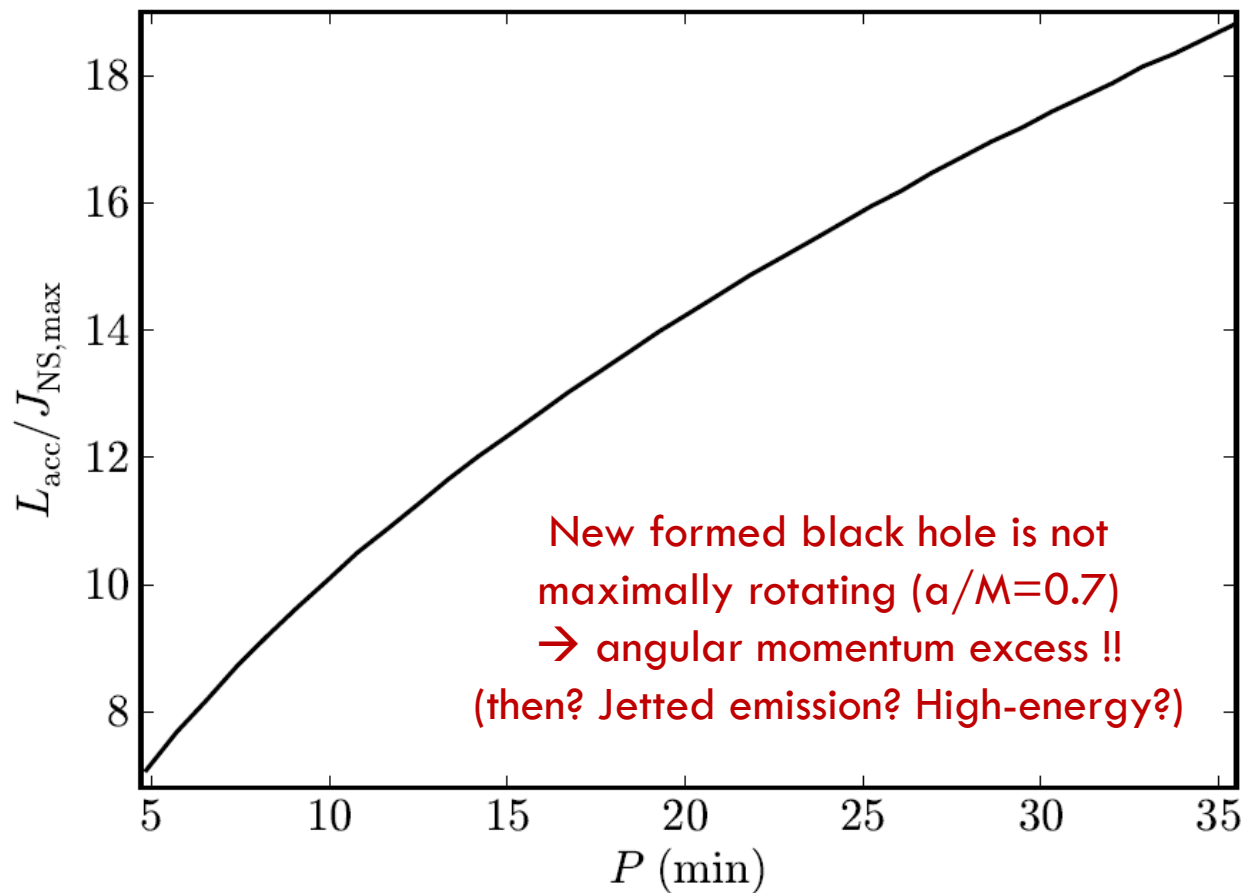


NS evolution up to the instability point



On the role of angular momentum in BdHNe

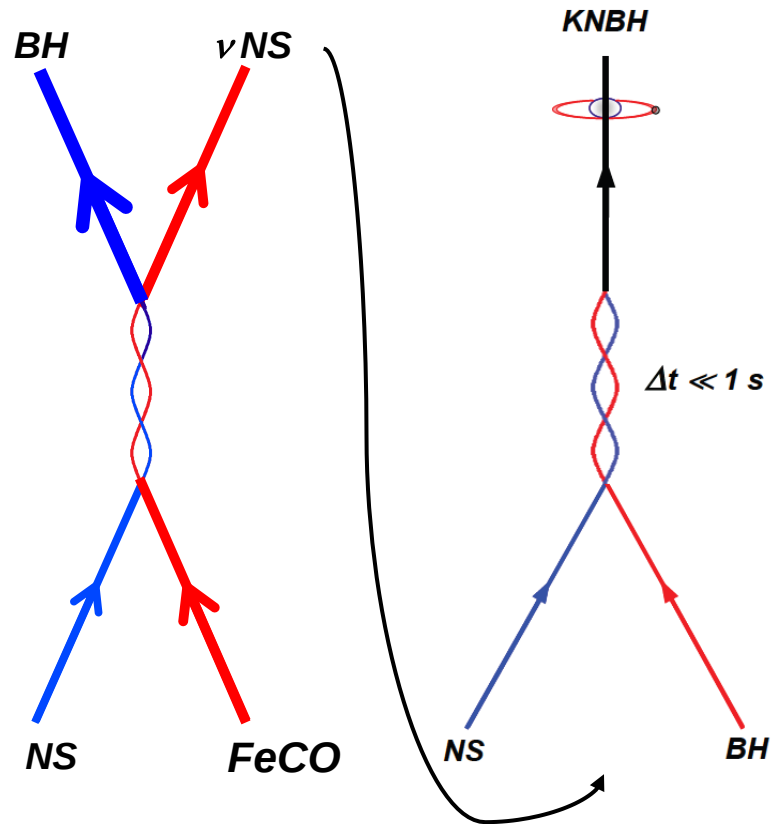
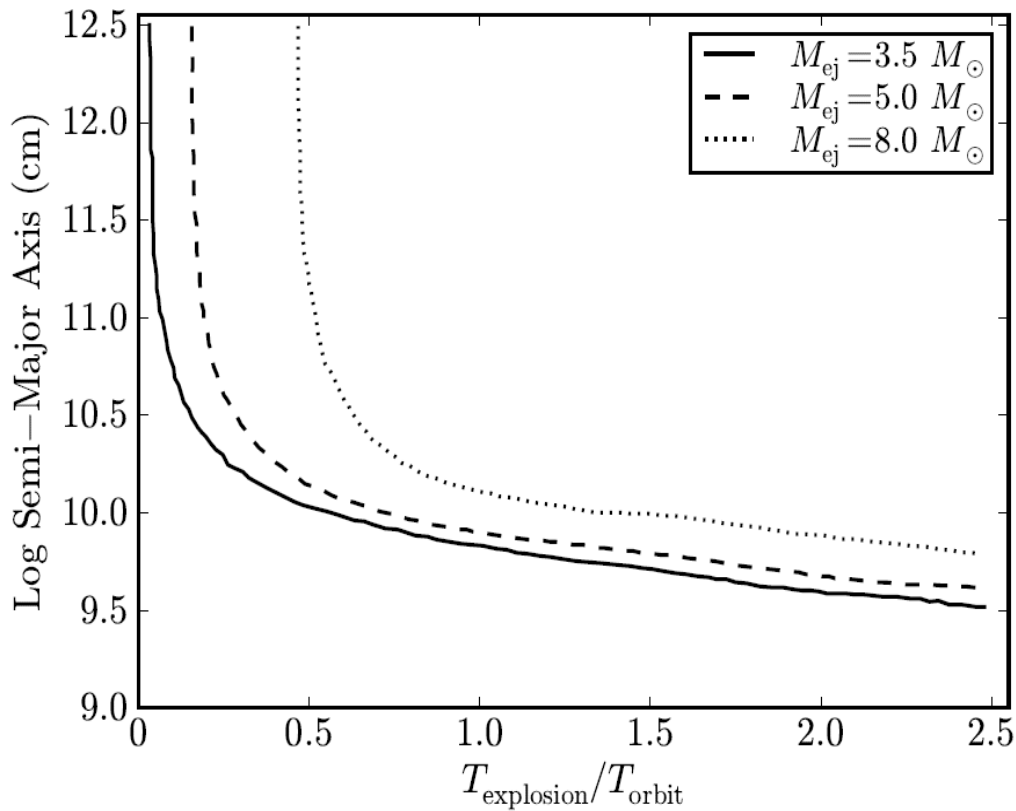
(Becerra, Cipolletta, Fryer, Rueda, Ruffini, ApJ 2015; arXiv: 1505.07580)



What's next?

On the NS-BH binaries produced by BdHNe

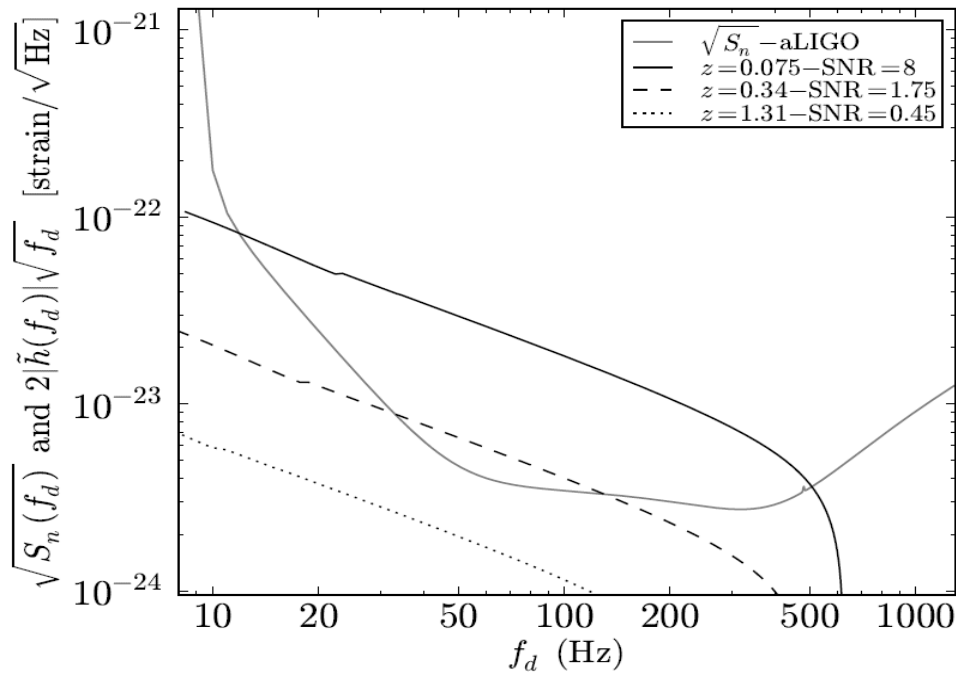
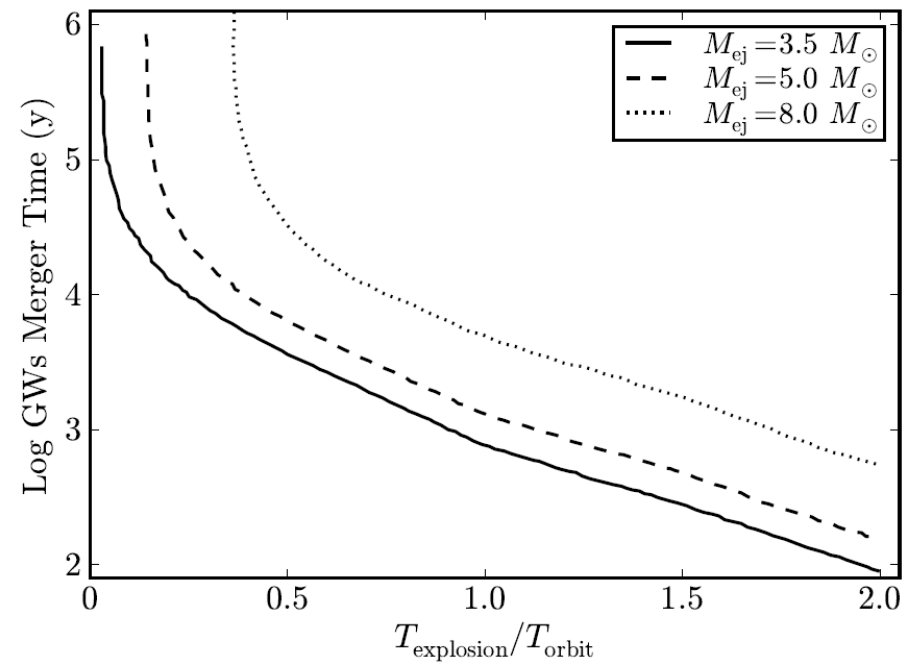
(Fryer, Oliveira, Rueda, Ruffini, Phys. Rev. Lett., in press)



What's next?

On the NS-BH binaries produced by BdHNe

(Fryer, Oliveira, Rueda, Ruffini, Phys. Rev. Lett., in press)



CONCLUSION

